

Article

Diffusion on Middle- ζ Cantor Sets

Alireza Khalili Golmankhaneh ^{1,*} , Arran Fernandez ²  and Ali Khalili Golmankhaneh ³
and Dumitru Baleanu ^{4,5}

¹ Department of Physics, Urmia Branch, Islamic Azad University, Urmia, Iran

² Department of Applied Mathematics and Theoretical Physics, University of Cambridge, Wilberforce Road, Cambridge CB3 0WA, UK; af454@cam.ac.uk

³ Young Researchers and Elite Club, Urmia Branch, Islamic Azad University, Urmia, Iran; a.khalili@tabrizu.ac.ir

⁴ Department of Mathematics, Cankaya University, Ankara 06530, Turkey; dimitru@cankaya.edu.tr

⁵ Institute of Space Sciences, P.O. Box, MG-23, R 76900 Magurele-Bucharest, Romania

* Correspondence: a.khalili@iaurmia.ac.ir; Tel.: +98-443-272-2702

Received: 23 May 2018; Accepted: 6 June 2018; Published: 2 July 2018



Abstract: In this paper, we study C^ζ -calculus on generalized Cantor sets, which have self-similar properties and fractional dimensions that exceed their topological dimensions. Functions with fractal support are not differentiable or integrable in terms of standard calculus, so we must involve local fractional derivatives. We have generalized the C^ζ -calculus on the generalized Cantor sets known as middle- ζ Cantor sets. We have suggested a calculus on the middle- ζ Cantor sets for different values of ζ with $0 < \zeta < 1$. Differential equations on the middle- ζ Cantor sets have been solved, and we have presented the results using illustrative examples. The conditions for super-, normal, and sub-diffusion on fractal sets are given.

Keywords: Hausdorff dimension; middle- ζ Cantor sets; staircase function; C^ζ -calculus; diffusion on fractal; random walk

1. Introduction

It is well known that many phenomena in nature can be modeled by fractals; these shapes can be observed almost anywhere in the natural world [1]. Fractal antennas have maximal length, area, and volume to accommodate a multi-band or wide-band design, which is useful in cellular telephone and microwave communications [2–4]. Fractals also play important roles in biology. For example, in the neural and vascular networks of the human body, viruses and certain tumors grow and ramify in a fractal shape [5–7]. In these studies, researchers tried to predict and recognize osteoporosis from test results and from the fractal structure of bone texture [8]. Fractals have also been hypothesized to be important for human perception of beauty in artworks [9,10]. Non-Markovian random walks and fractal dimensions which are connected to physical properties of fractal sets were studied in [11–13]. The polynomial asymptotic behavior of the Wiener index on infinite lattices including fractals has been given [14].

Anomalous diffusion on fractals has received attention in recent years from various researchers [15,16]. Non-constant diffusion coefficients have been studied [17], and the diffusion coefficient is proportional to a power of the noise intensity [18]. The models for the diffusion coefficient characterize random motions in various regions of parameterized space and relate to the fractal structure as a function of the slope of the map [19,20]. Fractal structural parameters have been applied to obtain porosity and tortuosity for micro-porous solids [21]. The scale-dependent fractal dimension for a random walk trajectory was used to derive the diffusion coefficient [22,23]. The quenched-trap model on a fractal lattice does not lead to continuous-time random walks if the spectral dimension is less than 2 [24].

Fractional calculus has been applied to define derivatives on fractal curves [25–32]. Fractional derivatives have non-local properties, so that they are used to model processes with memory effects [29,30].

Local fractional derivatives were suggested and applied from a physics perspective [33,34] in a formalism called C^ζ -calculus (or C^ζ -C), which has also been generalized for unbounded and singular functions [35]. Schrödinger equations on fractal curves were derived using C^ζ -C and Feynman path methods [36]. A mathematical model of diffraction was given for fractal sets [37]. Non-local derivatives were defined for fractal sets and applied in fractal mediums [38–40]. The Fokker–Planck equation for thick fractal absorbers was derived in view of C^ζ -C [41].

Anomalous diffusion in fractal comb structures, where the anomalous diffusion exponent depends on the fractal dimension of the comb structure, has also been studied [42–44], and the theoretical predictions in these works was recently experimentally proven [45]. Recently, as an application of the mathematical model, experimental and simulation results were utilized to model sub-diffusion and super-diffusion in physical processes [46,47].

By conducting research along these lines, we have generalized C^ζ -C to middle- ζ Cantor sets.

The outline of the paper is as follows. In Section 2, we review C^ζ -C and the basic tools required. In Section 3, we apply the C^ζ -C on the middle- ζ Cantor sets. In Section 4, we consider and solve some differential equations on middle- ζ Cantor sets, and, in Section 5, we consider diffusion processes on such sets. Section 4 is devoted to the conclusion.

2. Basic Tools in the Fractal Calculus

In this section, we review the Cantor-like sets and their properties [48], and then summarize C^ζ -C [33,34,38,39].

2.1. Middle- ζ Cantor Sets

Let us consider a unit interval $J = [0, 1]$, and construct the middle- ζ Cantor fractal set C^ζ from it as follows.

In the first step, we remove an open interval of length ζ from the exact middle of the interval I , to obtain:

$$C_1^\zeta = \left[0, \frac{1-\zeta}{2}\right] \cup \left[\frac{1+\zeta}{2}, 1\right]. \quad (1)$$

In the second step, we pick up two open disjoint intervals with length ζ^2 from the middle of each of the remaining intervals that comprise the set C_1^ζ , in order to obtain

$$C_2^\zeta = \left[0, \frac{1-\zeta-2\zeta^2}{4}\right] \cup \left[\frac{1-\zeta+2\zeta^2}{4}, \frac{1-\zeta}{2}\right] \cup \left[\frac{1+\zeta}{2}, \frac{3+\zeta-2\zeta^2}{4}\right] \cup \left[\frac{3+\zeta+2\zeta^2}{4}, 1\right]. \quad (2)$$

After iterating this process infinitely many times, with the set constructed at stage k being denoted by C_k^ζ , we obtain the definition of the middle- ζ Cantor set as follows:

$$C^\zeta = \bigcap_{k=1}^{\infty} C_k^\zeta. \quad (3)$$

It is clear that the set C^ζ has a self-similarity property, which makes it easy for us to find its fractional dimension. Namely, for every middle- ζ Cantor set, the Hausdorff dimension is given by

$$\dim_H(C^\zeta) = \frac{\log 2}{\log(1-\zeta)}, \quad (4)$$

where $H(C^\xi)$ is the Hausdorff measure that was used to derive Hausdorff dimension [48]. Furthermore, the middle- ξ Cantor set has zero Lebesgue measure because [48]:

$$L_m(C^\xi) = \lim_{k \rightarrow \infty} L_m(C_k^\xi) = \lim_{k \rightarrow \infty} (1 - \xi)^k = 0. \quad (5)$$

Remark 1. If we choose $\xi = 1/3$, $\xi = 1/4$, $\xi = 1/5$, $\xi = 3/5$, then we obtain the Cantor triadic set, 4-adic-type Cantor-like set, and 5-adic-type Cantor-like set, respectively. Note that middle- ξ Cantor sets with higher ξ have a larger measure; this was shown by using the idea of differences between Cantor sets [49]. See Section 3 below for more details on these sets.

2.2. Local Fractal Calculus

If C^ξ is a middle- ξ Cantor set contained in an interval $J = [v, w] \subset \mathbb{R}$, then the flag function for C^ξ is indicated by $\varphi(C^\xi, J)$ and defined by [33,34]

$$\varphi(C^\xi, J) = \begin{cases} 1, & \text{if } C^\xi \cap J \neq \emptyset, \\ 0, & \text{otherwise.} \end{cases} \quad (6)$$

For a set C^ξ and a subdivision $Q_{[v,w]} = \{v = y_0, y_1, y_2, \dots, y_n = w\}$ of the interval $J = [v, w]$, we define

$$\rho^\xi[C^\xi, J] = \sum_{i=1}^n \Gamma(\xi + 1)(y_i - y_{i-1})^\xi \varphi(C^\xi, [y_{i-1}, y_i]) \quad (7)$$

for any ξ with $0 < \xi \leq 1$. Given $\delta > 0$, the associated coarse-grained mass function $\gamma_\delta^\xi(C^\xi, v, w)$ of the intersection $C^\xi \cap [v, w]$ is given by

$$\gamma_\delta^\xi(C^\xi, v, w) = \inf_{Q_{[v,w]}: |Q| \leq \delta} \rho^\xi[C^\xi, J], \quad (8)$$

where the infimum is taken over all subdivisions Q of $[v, w]$ satisfying $|Q| := \max_{1 \leq i \leq n} (y_i - y_{i-1}) \leq \delta$. Then, the mass function $\gamma^\xi(C^\xi, v, w)$ is given by [33,34]

$$\gamma^\xi(C^\xi, v, w) = \lim_{\delta \rightarrow 0} \gamma_\delta^\xi(C^\xi, v, w). \quad (9)$$

The integral staircase function $S_{C^\xi}^\xi(y)$ of order ξ for a fractal set C^ξ is defined in [33,34] by

$$S_{C^\xi}^\xi(y) = \begin{cases} \gamma^\xi(C^\xi, v_0, y), & \text{if } y \geq a_0, \\ -\gamma^\xi(C^\xi, v_0, y), & \text{otherwise,} \end{cases} \quad (10)$$

where v_0 is an arbitrary real number. A point y is a point of change of a function $u(y)$ that is not constant over any open interval (v, w) involving y . All points of change of y is named the set of change of $u(y)$ and is indicated by $\text{Sch}(S_{C^\xi}^\xi)$ [33,34]. If $\text{Sch}(S_{C^\xi}^\xi)$ is a closed set and every point in it is a limit point, then $\text{Sch}(S_{C^\xi}^\xi)$ is called ξ -perfect.

The ξ -dimension of $C^\xi \cap [v, w]$ is

$$\begin{aligned} \dim_\xi(C^\xi \cap [v, w]) &= \inf\{\xi : \gamma^\xi(C^\xi, v, w) = 0\} \\ &= \sup\{\xi : \gamma^\xi(C^\xi, v, w) = \infty\}. \end{aligned} \quad (11)$$

We also define the concepts of C^ξ -limits and C^ξ -continuity, which will be used in the next section.

For a function $h : C^\zeta \rightarrow \mathbb{R}$ and a point $x \in C^\zeta$, a number l is said to be the limit of h through the points of C^ζ , or simply the C^ζ -limit of h as $z \rightarrow x$, if given any $\epsilon > 0$ there exists $\delta > 0$ such that

$$z \in C^\zeta \quad \text{and} \quad |z - x| < \delta \Rightarrow |h(z) - l| < \epsilon. \quad (12)$$

If such a number exists, then it is denoted by:

$$l = C^\zeta\text{-}\lim_{z \rightarrow x} h(z). \quad (13)$$

A function $h : C^\zeta \rightarrow \mathbb{R}$ is said to be C^ζ -continuous at $x \in C^\zeta$ if

$$h(x) = C^\zeta\text{-}\lim_{z \rightarrow x} h(z). \quad (14)$$

2.3. C^ζ -Differentiation

If C^ζ is an ζ -perfect set, then the C^ζ -derivative of a function u defined on C^ζ at a point y is defined to be the following, assuming the limit exists [33,34]:

$$D_{C^\zeta}^\zeta u(y) = \begin{cases} C^\zeta\text{-}\lim_{z \rightarrow y} \frac{f(z) - f(y)}{S_{C^\zeta}^\zeta(z) - S_{C^\zeta}^\zeta(y)}, & \text{if } z \in C^\zeta, \\ 0, & \text{otherwise.} \end{cases} \quad (15)$$

Let u be a bounded function on C^ζ and J be a closed interval as above [33,34]. Then, we define

$$\mathfrak{M}[u, C^\zeta, J] = \sup_{y \in C^\zeta \cap J} u(y) \quad \text{if } C^\zeta \cap J \neq \emptyset \quad (16)$$

$$= 0 \quad \text{otherwise,} \quad (17)$$

and similarly

$$\mathfrak{m}[u, C^\zeta, J] = \inf_{y \in C^\zeta \cap J} u(y) \quad \text{if } C^\zeta \cap J \neq \emptyset \quad (18)$$

$$= 0 \quad \text{otherwise.} \quad (19)$$

If $S_{C^\zeta}^\zeta(y)$ is finite for $y \in [v, w]$, and $Q = \{v = y_0, y_1, \dots, y_n = w\}$ is a subdivision of $[v, w]$, then the upper C^ζ -sum and lower C^ζ -sum for a function u over the subdivision Q are given, respectively, by [33,34]

$$\mathfrak{U}^\zeta[u, C^\zeta, Q] = \sum_{j=1}^m \mathfrak{M}[u, C^\zeta, [y_j, y_{j-1}]] (S_{C^\zeta}^\zeta(y_j) - S_{C^\zeta}^\zeta(y_{j-1})) \quad (20)$$

and

$$\mathfrak{L}^\zeta[u, C^\zeta, Q] = \sum_{j=1}^m \mathfrak{m}[u, C^\zeta, [y_j, y_{j-1}]] (S_{C^\zeta}^\zeta(y_j) - S_{C^\zeta}^\zeta(y_{j-1})). \quad (21)$$

Let u be a bounded function on C^ζ . We say that u is C^ζ -integrable on $[v, w]$ if [33,34] the two quantities

$$\int_v^w u(y) d_{C^\zeta}^\zeta y = \sup_{Q[v,w]} \mathfrak{U}^\zeta[u, C^\zeta, Q], \quad (22)$$

$$\int_v^w u(y) d_{C^\zeta}^\zeta y = \inf_{Q[v,w]} \mathfrak{L}^\zeta[u, C^\zeta, Q], \quad (23)$$

are equal. In that case, the C^ζ -integral of u on $[v, w]$ is denoted by $\int_v^w u(y) d_{C^\zeta}^\zeta y$ and is given by the common value of Labels (22)–(23).

Fundamental Theorem of C^ζ -Calculus. Suppose that $u(y) : C^\zeta \rightarrow \mathbb{R}$ is C^ζ -continuous and bounded on C^ζ . If we define $g(z)$ by

$$g(z) = \int_a^z f(y) d_{C^\zeta}^\zeta y \quad (24)$$

for all $y \in [v, w]$, then:

$$D_{C^\zeta}^\zeta g(y) = f(y) \chi_{C^\zeta}(y), \quad (25)$$

where $\chi_{C^\zeta}(y)$ is the characteristic function of the middle- ζ Cantor set.

Conversely, if $f(y)$ is an C^ζ -differentiable function, then we have [33,34]

$$D_{C^\zeta}^\zeta f(y) = h(y) \chi_{C^\zeta}(y) \quad (26)$$

for some function h , and, consequently, it follows that

$$\int_v^w h(y) d_{C^\zeta}^\zeta y = f(b) - f(a). \quad (27)$$

3. Staircase Functions on Middle- ζ Cantor Sets

In this section, we plot middle- ζ Cantor sets and their staircase functions $S_{C^\zeta}^\zeta(y)$ for special cases, in order to present details of the paper.

3.1. The Cantor Triadic Set

The Cantor triadic set is generated by iteration as follows:

- *Step 1.* Remove an open interval of length $1/3$ from the middle of the interval $J = [0, 1]$.
- *Step 2.* Remove an open interval of length $(1/3)^2$ from the middle of each one of the closed intervals with length $1/3$ remaining from step 1.
- ...
- *Step k .* Remove an open interval of length $(1/3)^k$ from the middle of each one of the closed intervals with length $(1/3)^{k-1}$ remaining from step $k - 1$.

In the case of the Cantor triadic sets, utilizing Equations (4) and (11), we get ζ -dimension as follows:

$$\dim_\zeta(C^{1/3} \cap [v, w]) = \dim_H(C^{1/3}) = 0.63. \quad (28)$$

In Figure 1a, we draw the process mentioned above that established the Cantor triadic set.

Using Equation (10), the staircase function of the Cantor triadic set ($S_{C^{1/3}}^{0.63}(x)$) is sketched in Figure 1c.

The staircase functions have important roles in C^ζ -C, being used in integration and differentiation of functions with fractal support.

3.2. The 5-Adic-Type Cantor-Like Set

The procedure to achieve the 5-adic-type Cantor-like set is similar to that for the Cantor triadic set, only differed to remove the open interval of length $3/5$ in every stage. We exhibit these steps in Figure 1b.

The ζ -dimension of the 5-adic-type Cantor-like set, considering Equations (4) and (11), is

$$\dim_\zeta(C^{3/5} \cap [v, w]) = \dim_H(C^{3/5}) = 0.43. \quad (29)$$

From Equation (10), the staircase function of the 5-adic-type Cantor-like ($S_{C^{3/5}}^{0.43}(y)$) is plotted in Figure 1d.

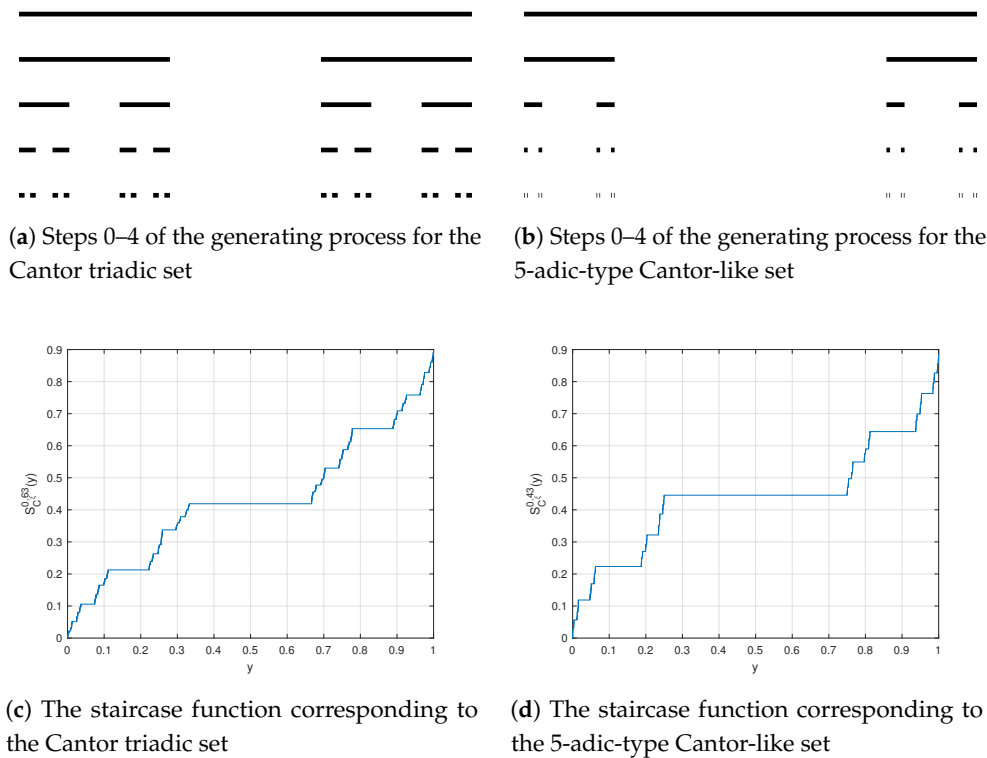


Figure 1. Basic properties of some example Cantor sets.

4. Differential Equations on Middle- ζ Cantor Sets

In this section, first, we study the integration and differentiation of functions whose support is a middle- ζ Cantor set. Secondly, differential equations formulated on middle- ζ Cantor sets are suggested and solved using illustrative examples.

Example 1. Consider a function with the fractal Cantor triadic set support as follows:

$$f(x) = \sin \left(2\pi x \chi_{C^{1/3}}^{0.63}(x) \right) \quad (30)$$

where

$$\chi_{C^{1/3}}^{0.63}(x) = \begin{cases} \frac{1}{\Gamma(1+0.63)}, & x \in C^{1/3}, \\ 0, & \text{otherwise,} \end{cases} \quad (31)$$

is the characteristic function of the fractal Cantor triadic set. We plot the function $f(x)$ in Figure 2a.

The C^ζ -derivative of $f(x)$, using conjugacy of C^ζ -C and ordinary calculus [33,34], is as follows:

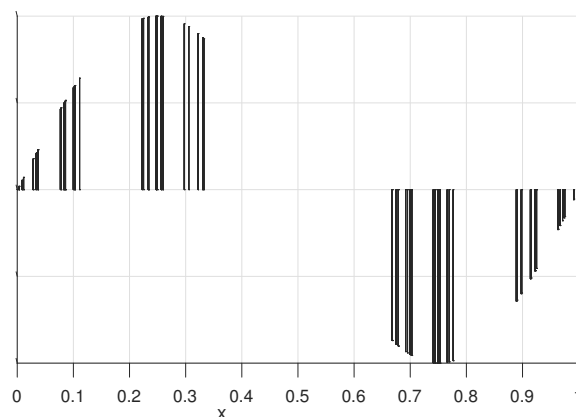
$$D_{C^{1/3}}^{0.63} f(x) = \frac{2\pi}{\Gamma(1+0.63)} \cos \left(2\pi x \chi_{C^{1/3}}^{0.63}(x) \right), \quad (32)$$

where $\Gamma(\cdot)$ denotes the gamma function. A plot of the function $D_{C^{1/3}}^{0.63} f(x)$ is shown in Figure 2b.

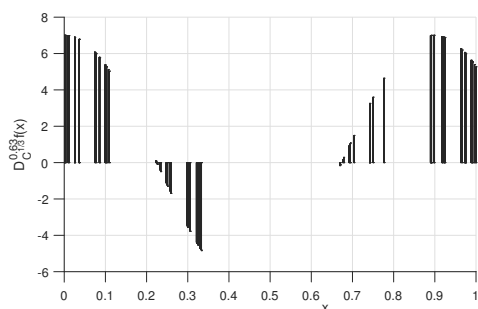
The C^ζ -integration of $f(x)$, considering conjugacy of C^ζ -C between the ordinary calculus [33,34], will be as follows:

$$\begin{aligned}
 \int_0^1 \sin \left(2\pi x \chi_{C^{1/3}}^{0.63}(x) \right) d_{C^{1/3}}^{0.63} x &= \frac{-\Gamma(1+0.63)}{2\pi} \left[\cos \left(2\pi \frac{S_{C^{1/3}}^{0.63}(x)}{\Gamma(1+0.63)} \right) \right]_0^1 \\
 &= \frac{1}{2\pi} \left[1 - \cos \left(2\pi \frac{S_{C^{1/3}}^{0.63}(1)}{\Gamma(1+0.63)} \right) \right] \\
 &= 0,
 \end{aligned}$$

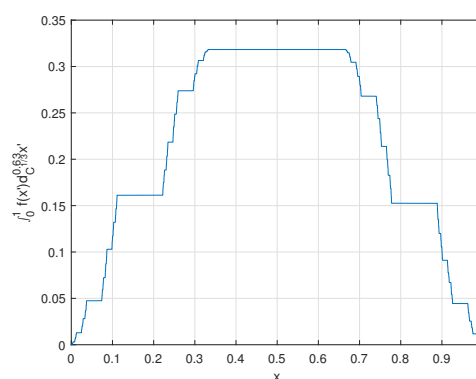
where we use $S_{C^{1/3}}^{0.63}(1) = \Gamma(1+0.63)$. In Figure 2c, we plot the integral function of $f(x)$ over $[0, 1]$.



(a) Graph of the function $f(x)$



(b) Graph of the fractal derivative $D_{C^{1/3}}^{0.63} f(x)$



(c) Graph of the fractal integral $\int_0^x f(x') d_{C^{1/3}}^{0.63} x'$

Figure 2. Graphs relevant to Example 1.

Example 2. Suppose we have a function on the fractal 5-adic-type Cantor-like set as follows:

$$g(x) = x^2 \chi_{C^{3/5}}^{0.43}(x). \quad (33)$$

This function $g(x)$ is sketched in Figure 3a.

The C^{ξ} -derivative of $g(x)$ is derived by a similar method as used in Example 1, which yields the following result:

$$D_{C^{3/5}}^{0.43} g(x) = 2x \chi_{C^{3/5}}^{0.43}(x). \quad (34)$$

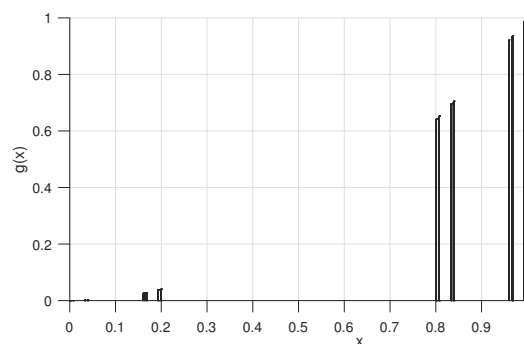
We plot $D_{C^{3/5}}^{0.43} g(x)$ in Figure 3b.

In the same manner, we obtain the F^ζ -integral of $g(x)$ as follows:

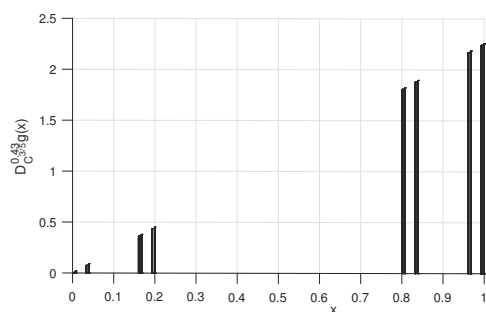
$$\begin{aligned} \int_0^1 (\chi_{C^{3/5}}^{0.43}) x^2 d_{C^{3/5}}^{0.43} x &= \frac{1}{3\Gamma(1+0.43)} \left(S_{C^{3/5}}^{0.43}(x) \right)^3 \Big|_0^1 \\ &= \frac{1}{3\Gamma(1+0.43)} \left(S_{C^{3/5}}^{0.43}(1) \right)^3 = 0.26. \end{aligned} \quad (35)$$

We plot the C^ζ -integral of $g(x)$ in Figure 3c.

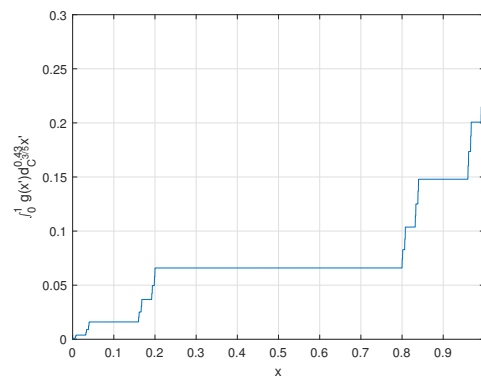
We are going to use the above results and examples for solving differential equations on the middle- ζ Cantor sets.



(a) Graph of the function $g(x)$



(b) Graph of the fractal derivative $D_{C^{3/5}}^{0.43} g(x)$



(c) Graph of the fractal integral $\int_0^x g(x') d_{C^{3/5}}^{0.43} x'$

Figure 3. Graphs relevant to Example 2.

5. Diffusion on Middle- ζ Cantor Sets

In this section, we define and give conditions for super, normal and sub-diffusion on middle- ζ Cantor sets.

5.1. Super-Diffusion

Let us consider time as continuous and space as a middle- ζ Cantor set. We consider a probability function $W(x, t)$, which is an C^ζ -differentiable function of the space coordinate x and is a differentiable

function of time t in the sense of standard calculus. The fractal diffusion equation for a random walk is suggested with the conditional probability $W(x, t)$ as follows:

$$\chi_{C^\zeta}(x) \frac{\partial W(x, t)}{\partial t} = K_{C^\zeta} (D_{C^\zeta, x}^\zeta)^2 W(x, t), \quad t \in \mathbb{R}, x \in C^\zeta, \quad (36)$$

with the initial condition

$$W(x, t = 0) = \delta_{C^\zeta}^\zeta(x), \quad (37)$$

where $[K_{C^\zeta}] = (Length^{2\zeta}/Time)$ is a generalized diffusion coefficient, $\delta_{C^\zeta}^\zeta(x)$ is the Dirac delta function with fractal support. Using conjugacy of C^ζ -C between standard calculus [33,34], we have the solution for Equation (36) as follows:

$$W(x, t) = \frac{t^{-1/2}}{\sqrt{4\pi K_{C^\zeta}}} \exp \left[\frac{-S_{C^\zeta}^\zeta(x)^2}{4K_{C^\zeta}t} \right]. \quad (38)$$

Since $S_{C^\zeta}^\zeta(x) \leq x^\zeta$, then Equation (38) can be written as:

$$W(x, t) \mapsto \frac{t^{-1/2}}{\sqrt{4\pi K_{C^\zeta}}} \exp \left[\frac{-x^{2\zeta}}{4K_{C^\zeta}t} \right]. \quad (39)$$

The function $W(x, t)$ is indicated as the probability distribution of super-diffusion on Cantor sets. Accordingly, the mean square random walk is

$$\langle S_{C^\zeta}^\zeta(x)^2 \rangle = 4K_{C^\zeta}t. \quad (40)$$

Using the upper bound $S_{C^\zeta}^\zeta(x)^2 \leq x^{2\zeta}$, we have

$$\langle x^2 \rangle \mapsto 4K_{C^\zeta}t^{1/\zeta}. \quad (41)$$

5.2. Normal Diffusion

Let us consider space as a middle- ζ Cantor set and fractal time associated with the middle- ζ Cantor set, both sets having the same value of ζ and the same dimension ζ . The random walk conditional probability $W(x, t)$ is given by

$$D_{C^\zeta, t}^\zeta W(x, t) = G_{C^\zeta} (D_{C^\zeta, x}^\zeta)^2 W(x, t), \quad (42)$$

where $[G_{C^\zeta}] = (Length^{2\zeta}/Time^\zeta)$ is a diffusion coefficient. The solution for Equation (42) with the initial condition Equation (37), utilizing conjugacy of C^ζ -C between standard calculus, is:

$$W(x, t) = \frac{S_{C^\zeta}^\zeta(t)^{-1/2}}{\sqrt{4\pi G_{C^\zeta}}} \exp \left[\frac{-S_{C^\zeta}^\zeta(x)^2}{4G_{C^\zeta}S_{C^\zeta}^\zeta(t)} \right]. \quad (43)$$

Considering the upper bound on the $S_{C^\zeta}^\zeta(\cdot)$, we obtain:

$$W(x, t) \mapsto \frac{t^{-\zeta/2}}{\sqrt{4\pi G_{C^\zeta}}} \exp \left[\frac{-x^{2\zeta}}{4G_{C^\zeta}t^\zeta} \right]. \quad (44)$$

The function $W(x, t)$ indicates the probability distribution for normal diffusion with a non-Gaussian propagator.

Applying conjugacy of C^ζ -C between standard calculus, we arrive at the mean square of displacement

$$\langle S_{C^\zeta}^\zeta(x)^2 \rangle = 4G_{C^\zeta} S_{C^\zeta}^\zeta(t), \quad (45)$$

and if we use the upper bound on $S_{C^\zeta}^\zeta(\cdot)$, we can write Equation (45) as follows:

$$\langle x^2 \rangle \mapsto 4G_{C^\zeta} t. \quad (46)$$

5.3. Sub-Diffusion

Let us consider time as a middle- ζ Cantor set with dimension β and space as a middle- ζ Cantor set with dimension ζ . A random walk on this fractal space-time has conditional probability that can be obtained by the following differential equation:

$$D_{C^{\zeta'}, t}^\beta W(x, t) = \chi_{C^{\zeta'}, t} L_{C^\zeta} (D_{C^\zeta, x}^\zeta)^2 W(x, t), \quad (47)$$

where $[L_{C^\zeta}] = (\text{Length}^{2\zeta} / \text{Time}^\beta)$ is a diffusion coefficient. Solving Equation (42) with the initial condition Equation (37), using conjugacy of C^ζ -C between standard calculus, one can obtain

$$W(x, t) = \frac{S_{C^\zeta}^\beta(t)^{-1/2}}{\sqrt{4\pi L_{C^\zeta}}} \exp \left[\frac{-S_{C^\zeta}^\zeta(x)^2}{4L_{C^\zeta} S_{C^\zeta}^\beta(t)} \right]. \quad (48)$$

In view of the upper bounds on $S_{C^\zeta}^\zeta(\cdot)$, we get

$$W(x, t) \mapsto \frac{t^{-\beta/2}}{\sqrt{4\pi L_{C^\zeta}}} \exp \left[\frac{-x^{2\zeta}}{4L_{C^\zeta} t^\beta} \right]. \quad (49)$$

The function $W(x, t)$ is named as the probability of sub-diffusion for a random walk as indicated above.

Similarly to the previous cases, the mean square of displacement in this case will be

$$\langle S_{C^\zeta}^\zeta(x)^2 \rangle = 4L_{C^\zeta} S_{C^\zeta}^\beta(t), \quad (50)$$

and in the same manner we use upper bounds on $S_{C^\zeta}^\zeta(\cdot)$ to get

$$\langle x^2 \rangle \mapsto 4L_{C^\zeta} t^{\beta/\zeta}. \quad (51)$$

Example 3. Consider a random walk model on the fractal 3-adic-type Cantor-like set. The corresponding mean square value displacement of the random walk is given by:

$$\langle S_{C^\zeta}^{0.63}(x)^2 \rangle = 4L_{C^\zeta} S_{C^\zeta}^\beta(t), \quad (52)$$

or

$$\langle x^2 \rangle \mapsto 4L_{C^\zeta} t^{\beta/0.63}, \quad (53)$$

where the respective cases $\beta > 0.63$, $\beta < 0.63$, and $\beta = 0.63$ are called super-diffusion, sub-diffusion and normal diffusion on the fractal 3-adic-type Cantor-like set, respectively.

In Figure 4, we draw the graphs of mean square value displacement random walk model for super-diffusion, sub-diffusion, and normal diffusion on fractal 3-adic-type Cantor-like sets in the case $\zeta = 0.63$.

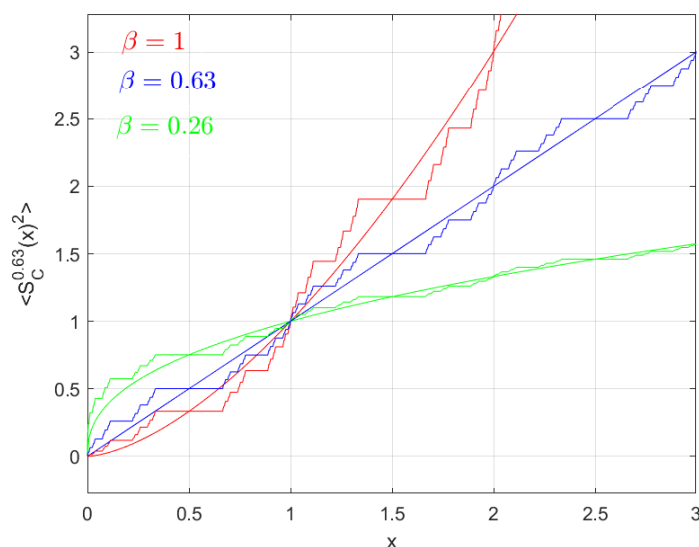


Figure 4. Graphs of the mean value displacement of the random walk model for super-diffusion (in red), sub-diffusion (in green) and normal diffusion (in blue) in the case $\alpha = 0.63$; the smoothed curves represent the corresponding processes on the real line.

Remark 2. We conclude the results of Section 5 as follows:

1. The diffusion is super-diffusion on the middle- ζ Cantor set if $\zeta < \beta$.
2. The diffusion is normal on the middle- ζ Cantor set if $\zeta = \beta$.
3. The diffusion is sub-diffusion on the middle- ζ Cantor set if $\zeta > \beta$.

Remark 3. In some figures, we have plotted bars instead of points for the graphs of functions with fractal support, in order to make the results more clear.

6. Conclusions

The C^ζ -calculus is a generalization of ordinary calculus that can be applied on middle- ζ Cantor sets for different values of ζ . Functions with middle- ζ Cantor set support were considered, and their derivatives and integrals were derived using C^ζ -calculus, which shows the advantage of using C^ζ -calculus over standard calculus. C^ζ -derivatives on new fractal sets were discussed and compared for functions with different fractal supports. New differential equations involving C^ζ -derivatives on middle- ζ Cantor sets have been suggested, which can be used as mathematical models for many physical problems. For example, we suggest conditions for super-, normal, and sub-diffusion on fractal sets.

Author Contributions: All authors contributed equally to this work. All authors have read and approved the final manuscript.

Funding: The second author is on a PhD studentship which was funded by the Engineering and Physical Sciences Research Council, UK.

Conflicts of Interest: The authors declare no conflict of interest.

References

1. Barnsley, M.F. *Fractals Everywhere*; Academic Press: New York, NY, USA, 2014.
2. Cohen, N. *Fractal Antenna Applications in Wireless Telecommunications: Electronics Industries Forum of New England*; IEEE: Toulouse, France, 1997.
3. Cohen, N. Fractal Antennas: Part 2. *Commun. Q.* **1996**, *44*, 53–66.

4. Werner, D.H.; Haupt, R.; Werner, P.L. Fractal Antenna Engineering: The Theory and Design of Fractal Antenna Arrays. *IEEE Antennas Propag. Mag.* **1999**, *41*, 37–59. [[CrossRef](#)]
5. Gazit, Y.; Baish, J.W.; Sfabakhsh, N.; Leunig, M.; Baxter, L.T.; Jain, R.K. Fractal characteristics of tumor vascular architecture during tumor growth and regression. *Microcirculation* **1997**, *4*, 395–402. [[CrossRef](#)] [[PubMed](#)]
6. Gazit, Y.; Berk, D.A.; Leunig, M.; Baxter, L.T.; Jain, R.K. Scale-invariant behavior and vascular network formation in normal and tumor tissue. *Phys. Rev. Lett.* **1995**, *75*, 2428–2431. [[CrossRef](#)] [[PubMed](#)]
7. Baish, J.W.; Jain, R.K. Fractals and cancer. *Cancer Res.* **2002**, *60*, 3683–3688.
8. Koh, K.J.; Park, H.N.; Kim, K.A. Prediction of age-related osteoporosis using fractal analysis on panoramic radiographs. *Imaging Sci. Dent.* **2012**, *42*, 231–235. [[CrossRef](#)] [[PubMed](#)]
9. Taylor, R.P.; Micolich, A.P.; Jonas, D. Fractal analysis of Pollock’s drip paintings. *Nature* **1999**, *399*, 422. [[CrossRef](#)]
10. Bountis, T.; Fokas, A.S.; Psarakis, E.Z. Fractal analysis of tree paintings by Piet Mondrian (1872–1944). *Int. J. Arts Technol.* **2017**, *10*, 27–42. [[CrossRef](#)]
11. Balankin, A.S.; Mena, B.; Martínez-González, C.L.; Matamoros, D.M. Random walk in chemical space of Cantor dust as a paradigm of superdiffusion. *Phys. Rev. E* **2012**, *86*, 052101. [[CrossRef](#)] [[PubMed](#)]
12. Balankin, A.S. Effective degrees of freedom of a random walk on a fractal. *Phys. Rev. E* **2015**, *92*, 062146. [[CrossRef](#)] [[PubMed](#)]
13. Golmankhaneh, A.K.; Baleanu, D. New heat and Maxwell’s equations on Cantor cubes. *Rom. Rep. Phys.* **2017**, *69*, 109.
14. Ori, O.; Cataldo, F.; Vukicevic, D.; Graovac, A. Wiener way to dimensionality. *Iranian J. Math. Chem.* **2010**, *1*, 5–15.
15. Poirier, D.R.; Geiger, G.H. Fick’s Law and Diffusivity of Materials. In *Transport Phenomena in Materials Processing*; Springer: Cham, The Netherlands, 2016.
16. Ben-Avraham, D.; Havlin, S. *Diffusion and Reactions in Fractals and Disordered Systems*; Cambridge University Press: Cambridge, UK, 2000.
17. Petersen, J.S.; Mack, C.A.; Sturtevant, J.L.; Byers, J.D.; Miller, D.A. Nonconstant diffusion coefficients: Short description of modeling and comparison to experimental results. In *Advances in Resist Technology and Processing XII*; International Society for Optics and Photonics: Bellingham, WA, USA, 1995; Volume 2438, pp. 167–181.
18. Lindner, B. Diffusion Coefficient of a Brownian Particle with a Friction Function Given by a Power Law. *J. Stat. Phys.* **2008**, *130*, 523–533. [[CrossRef](#)]
19. Schell, M.; Fraser, S.; Kapral, R. Diffusive dynamics in systems with translational symmetry: A one-dimensional-map model. *Phys. Rev. A* **1982**, *26*, 504–521. [[CrossRef](#)]
20. Klages, R.; Dorfman, J.R. Simple Maps with Fractal Diffusion Coefficients. *Phys. Rev. Lett.* **1995**, *74*, 387–390. [[CrossRef](#)] [[PubMed](#)]
21. Schieferstein, E.; Heinrich, P. Diffusion Coefficients Calculated for Microporous Solids from Structural Parameters Evaluated by Fractal Geometry. *Langmuir* **1997**, *13*, 1723–1728. [[CrossRef](#)]
22. Gmachowski, L. Fractal model of anomalous diffusion. *Eur. Biophys. J.* **2015**, *44*, 613–621. [[CrossRef](#)] [[PubMed](#)]
23. Bujan-Núñez, M.C. Scaling behavior of Brownian motion interacting with an external field. *Mol. Phys.* **1998**, *94*, 361–371. [[CrossRef](#)]
24. Miyaguchi, T.; Akimoto, T. Anomalous diffusion in a quenched-trap model on fractal lattices. *Phys. Rev. E* **2015**, *91*, 010102. [[CrossRef](#)] [[PubMed](#)]
25. Uchaikin, V.V. *Fractional Derivatives for Physicists and Engineers Vol. 1 Background and Theory. Vol 2. Application*; Springer: Berlin, Germany, 2013.
26. Tatom, F.B. The relationship between fractional calculus and fractals. *Fractals* **1995**, *3*, 217–229. [[CrossRef](#)]
27. Kolwankar, K.M.; Gangal, A.D. Fractional differentiability of nowhere differentiable functions and dimensions. *Chaos* **1996**, *6*, 505–513. [[CrossRef](#)] [[PubMed](#)]
28. Nigmatullin, R.R.; Le Mehaute, A. Is there geometrical/physical meaning of the fractional integral with complex exponent. *J. Non-Cryst. Solids* **2005**, *351*, 2888–2899. [[CrossRef](#)]
29. Richard, H. *Fractional Calculus: An Introduction for Physicists*; World Scientific: London, UK, 2014.
30. Hilfer, R. (Ed.) *Applications of Fractional Calculus in Physics*; World Scientific: London, UK, 2000.

31. Asad, H.; Mughal, M.J.; Zubair, M.; Naqvi, Q.A. Electromagnetic Green's function for fractional space. *J. Electromagn. Wave* **2012**, *26*, 1903–1910. [[CrossRef](#)]
32. Zubair, M.; Mughal, M.J.; Naqvi, Q.A. *Electromagnetic Fields and Waves in Fractional Dimensional Space*; Springer: Berlin/Heidelberg, Germany, 2012.
33. Parvate, A.; Gangal, A.D. Calculus on fractal subsets of real-line I: Formulation. *Fractals* **2009**, *17*, 53–148. [[CrossRef](#)]
34. Parvate, A.; Gangal, A.D. Calculus on fractal subsets of real line II: Conjugacy with ordinary calculus. *Fractals* **2011**, *19*, 271–290. [[CrossRef](#)]
35. Golmankhaneh, A.K.; Baleanu, D. Fractal calculus involving Gauge function. *Commun. Nonlinear Sci.* **2016**, *37*, 125–130. [[CrossRef](#)]
36. Golmankhaneh, A.K.; Golmankhaneh, A.K.; Baleanu, D. About Schrödinger equation on fractals curves imbedding in R^3 . *Int. J. Theor. Phys.* **2015**, *54*, 1275–1282. [[CrossRef](#)]
37. Golmankhaneh, A.K.; Baleanu, D. Diffraction from fractal grating Cantor sets. *J. Mod. Opt.* **2016**, *63*, 1364–1369. [[CrossRef](#)]
38. Golmankhaneh, A.K.; Baleanu, D. New derivatives on the fractal subset of Real-line. *Entropy* **2016**, *18*, 1. [[CrossRef](#)]
39. Golmankhaneh, A.K.; Baleanu, D. Non-local Integrals and Derivatives on Fractal Sets with Applications. *Open Phys.* **2016**, *14*, 542–548. [[CrossRef](#)]
40. Golmankhaneh, A.K.; Tunc, C. On the Lipschitz condition in the fractal calculus. *Chaos Soliton Fract.* **2017**, *95*, 140–147. [[CrossRef](#)]
41. Ashrafi, S.; Golmankhaneh, A.K. Energy Straggling Function by F^α -Calculus. *ASME J. Comput. Nonlinear Dyn.* **2017**, *12*, 051010. [[CrossRef](#)]
42. Sandev, T.; Iomin, A.; Kantz, H. Anomalous diffusion on a fractal mesh. *Phys. Rev. E* **2017**, *95*, 052107. [[CrossRef](#)] [[PubMed](#)]
43. Sandev, T.; Iomin, A.; Méndez, V. Lévy processes on a generalized fractal comb. *J. Phys. A Math. Theor.* **2016**, *49*, 355001. [[CrossRef](#)]
44. Iomin, A. Subdiffusion on a fractal comb. *Phys. Rev. E* **2011**, *83*, 052106. [[CrossRef](#)] [[PubMed](#)]
45. Zhokh, A.; Trypolskyi, A.; Strizhak, P. Relationship between the anomalous diffusion and the fractal dimension of the environment. *Chem. Phys.* **2018**, *503*, 71–76. [[CrossRef](#)]
46. Golmankhaneh, A.K.; Balankin, A.S. Sub-and super-diffusion on Cantor sets: Beyond the paradox. *Phys. Lett. A* **2018**, *382*, 960–967. [[CrossRef](#)]
47. Balankin, A.S.; Golmankhaneh, A.K.; Patiño-Ortiz, J.; Patiño-Ortiz, M. Noteworthy fractal features and transport properties of Cantor tartans. *Phys. Lett. A* **2018**, *382*, 1534–1539. [[CrossRef](#)]
48. Robert, D.; Urbina, W. On Cantor-like sets and Cantor-Lebesgue singular functions. *arXiv* **2014**, arXiv:1403.6554. [[CrossRef](#)]
49. Kraft, R.L. What's the difference between Cantor sets? *Am. Math. Mon.* **1994**, *101*, 640–650. [[CrossRef](#)]

